

## The Crystallography of Acridine. Part III. The Structure of Acridine II

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A three-dimensional analysis of a second crystalline form of acridine is described. Acridine II is monoclinic with

$$a = 16.292, b = 18.831, c = 6.072 \text{ \AA}, \beta = 95^\circ 4', Z = 8$$

and space group  $P2_1/a$ . There are two molecules in the asymmetric unit. The structure has been determined by trial-and-error and refined by successive calculations of differential syntheses from the three-dimensional data. Termination-of-series errors have been corrected by Booth's back-shift method.

Determination of the best mean isotropic temperature factors for each atom has shown that one of the molecules in the asymmetric unit is subject to more intense apparent thermal motion than the other. It has also been shown that the atomic thermal motion in each molecule is anisotropic and that it increases with the distance of an atom from its molecular centre. This anisotropic thermal motion, which arises from rigid-body vibrations and librations of the molecules, has not been allowed for in detail in the present refinement.

The two molecules in the asymmetric unit exhibit significant and different departures from planarity, but at this stage of refinement there are no significant differences in bond lengths (r.m.s. deviation 0.011 Å). The weighted mean bond lengths (estimated standard deviation between 0.003 and 0.005 Å) agree well with those observed in acridine III (r.m.s. deviation 0.008 Å) and in anthracene (r.m.s. deviation 0.005 Å).

### Introduction

Acridine,  $C_{13}H_9N$ , is a heterocyclic molecule derived from anthracene by the replacement of one *meso*-CH group by nitrogen. Fig. 1 shows the structural formula with the atoms numbered in the order preferred by Albert (1951). The numbers in brackets are used to refer to a second molecule.

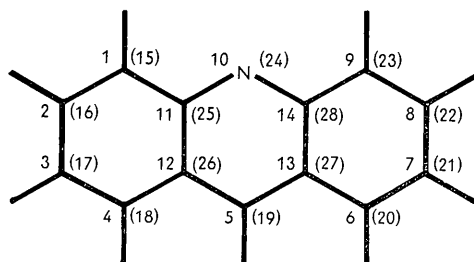


Fig. 1. Structural formula of acridine with numbering adopted for two molecules.

Three distinct crystalline forms of acridine have been described in Part I (Lowde, Phillips & Wood, 1953) and the structure of the simplest of these forms, acridine III, has been reported in detail in Part II (Phillips, 1956*a*). This paper describes a second of these forms, designated acridine II. This form is more complex crystallographically than acridine III, in

particular because the asymmetric unit in acridine II consists of two molecules.

The objects of the work now to be described were to determine the crystal structure of acridine II with sufficient accuracy to discover its relations with the other polymorphic forms of acridine, and to make further measurements of the molecular dimensions for comparison with those determined quite independently from acridine III. It is clear from the results obtained, however, that acridine II is a suitable structure in which to investigate directly the effects of molecular interactions on intramolecular dimensions. The two molecules in the asymmetric unit have markedly different environments and it is possible that sufficiently accurate measurements would reveal the different effects of these environments on them.

The possibility of measuring the effects of molecular interactions in crystals with more than one molecule per asymmetric unit has already been noted, for example by de Lange, Robertson & Woodward (1939) in their study of *trans*-azobenzene, but no thorough investigation has yet been made. There is, however, a growing body of evidence, for example from studies of crystal fluorescence (e.g. Sangster & Irvine, 1956), that molecular states are modified by interactions in the crystalline state. In particular it is known that 'fluorescence in aromatic compounds is associated with planarity' (Birks & Cameron, 1959) being determined by the presence and configuration of  $\pi$ -electrons in the molecule and that crystals of many normally

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planar molecules, including acridine, exhibit less than the expected fluorescence. In view of this evidence it is encouraging to find as the result of the present investigation that the two molecules in the asymmetric unit in acridine II exhibit significant and different departures from planarity. Their observed dimensions, on the other hand, are not significantly different at the present level of accuracy and the detailed investigation of the effects of molecular interactions on molecular dimensions, which are expected to be small, must await more accurate intensity measurements.

### Crystal data

The crystals are monoclinic prismatic (Lowde *et al.*, 1953). The cell dimensions have been re-determined by the Straumanis method using Cr  $K\alpha_1$  radiation,  $\lambda = 2.28962 \text{ \AA}$ , with the following results:

$$\begin{aligned} a &= 16.292 \pm 0.004, & b &= 18.831 \pm 0.002, \\ c &= 6.072 \pm 0.002 \text{ \AA}; & \beta &= 95^\circ 4' \pm 2'; \\ U &= 1855.6 \text{ \AA}^3; & Z &= 8. \end{aligned}$$

This gives the calculated density as  $1.283 \text{ g.ml.}^{-1}$  which may be compared with the measured value of  $1.278 \text{ g.ml.}^{-1}$ . Linear absorption coefficient, for X-rays with wavelength  $1.542 \text{ \AA}$ ;  $\mu = 6.92 \text{ cm.}^{-1}$ .

Reflexions  $h0l$  are absent when  $h$  is odd and reflexions  $0k0$  are absent when  $k$  is odd. No other reflexions are systematically absent. The space group thus is  $P2_1/a-C_{2h}^5$ .

### Intensity measurements

Roughly cylindrical crystals about  $0.5 \text{ mm.}$  in diameter were mounted on glass fibres with 'Duco' cement for rotation about the  $[001]$ ,  $[010]$  and  $[100]$  axes. The reflexions from these were recorded on equi-inclination Weissenberg photographs of the layers  $l = 0-5$ ,  $k = 0-7$  and  $h = 0-12$ . Each exposure was made with a set of three sheets of Kodak no-screen X-ray film. Exposure times varied up to 100 hr. All the estimations of reflexion intensity were made by visual comparison of the reflexion spots with those obtained by timed exposures of a single reflexion. Only the extended spots were measured on upper-level photographs and the intensities were corrected for extension (Phillips, 1956*b*). The Lorentz-polarization factors were derived from charts due to Lu (1943). No absorption corrections were applied.

The measurements from the various layer lines were brought to the same relative scale by the method described in Part II (Phillips, 1956*a*). In calculation of the average intensity for each reflexion the various estimations were weighted according to the inverse of the conversion factors required to bring them to the same scale. These conversion factors, after being adjusted to give a nearly absolute scale, varied from

1.72 to 5.32 with a mean value of 3.10. They were smallest for the densely populated  $[001]$ -axis photographs.

Within the Cu sphere of reflexion there are 4225 possible reflexions but only 2679, i.e. 63.4%, were estimated. This is a slightly smaller proportion than was observed for acridine III, but is about the same as that observed for anthracene. The intensities of 803 reflexions were estimated once, 1263 twice (i.e. on separate photographs taken about different axes) and 613 three times. A determination of the accuracy of these measurements made by the range-estimate method (Ibers, 1956) showed that  $\sigma(F)$  varies roughly as  $0.09F$  with a minimum value of 0.6. The final comparison of observed and calculated structure factors has shown, however, that the very largest amplitudes, which were estimated only once, are subject to rather larger errors. The observed structure amplitudes together with the final calculated values are listed in Table 1.

### Structure determination

The projection of the structure on  $(001)$  was considered first since it was expected to show the molecules resolved from one another and projected in much the same way as those in acridine III and anthracene. The Patterson projection was calculated and found to have narrow rows of peaks near the origin lying near the intersections with the  $\{420\}$  planes. These peaks were interpreted as arising from intramolecular vectors and showing the two molecules in the asymmetric unit in projection on  $(001)$  to be nearly parallel to  $(420)$ . Bragg-Lipson charts were prepared for the relatively strong reflexions 060, 080, 410, 420 and 430, and the weak one 440, and the positions of the molecules, projected as in acridine III, were adjusted to give good agreement for these structure factors. The structure factors for all  $hk0$  reflexions with  $\sin \theta < 0.4$  were then calculated and their signs were used in a calculation of  $\rho(xy)$ . In this first synthesis the individual atoms were not resolved but adjustments to the atomic positions derived from difference (error) syntheses in several cycles of refinement improved the map slowly until the atoms were resolved in one molecule ( $A$ ) but not in the other. A partial difference synthesis, calculated by taking out the contribution from the well-resolved molecule, then showed the second molecule resolved and clearly in a rather different orientation from that assumed. The two molecules in the asymmetric unit are rather differently inclined to  $(001)$  as may be seen in Fig. 4. Further refinement by the difference synthesis technique quickly reduced  $R$ , the residual error index (cf. Sutton, 1958, p. 27) of the  $hk0$  reflexions, to 0.23.

The  $z$ -coordinates also were determined by trial-and-error. Several models based on the type of packing found in acridine III were tested by the calculation of  $hk1$  structure factors and the one which gave satis-









Table 1 (cont.)

h	k	l	$10 F_0 $	$10\sigma_c$	h	k	l	$10 F_0 $	$10\sigma_c$	h	k	l	$10 F_0 $	$10\sigma_c$	h	k	l	$10 F_0 $	$10\sigma_c$	h	k	l	$10 F_0 $	$10\sigma_c$
16	00	03	195	166	17	00	00	23	42	17	07	01	11	5	18	01	00	13	22	18	07	00	11	19
16	00	02	165	150	17	00	02	69	58	17	07	02	11	9	18	01	01	24	22	18	07	01	11	19
16	00	01	166	151	17	00	03	14	35	17	07	03	27	28	18	01	02	76	60	18	07	02	11	19
16	00	00	149	126	17	00	04	47	74	17	07	04	7	5	18	01	03	30	30	18	07	03	11	19
16	00	00	9	10	17	00	05	13	2	17	07	05	143	116	18	01	04	19	20	18	07	04	11	19
16	10	00	24	24	17	00	06	33	2	17	07	06	15	30	18	01	05	45	45	18	07	05	11	19
16	10	01	147	83	17	00	07	21	2	17	07	07	142	142	18	01	06	39	39	18	07	06	11	19
16	10	02	36	2	17	00	08	13	10	17	07	08	31	26	18	01	07	31	32	18	07	07	11	19
16	10	03	16	13	17	00	09	21	12	17	07	09	28	30	18	01	08	83	64	18	07	08	11	19
16	10	04	126	13	17	00	10	12	8	17	07	10	40	40	18	01	09	83	7	18	07	09	11	19
16	10	05	88	110	17	00	11	98	74	17	07	11	38	18	18	01	10	9	9	18	07	10	11	19
16	10	06	3	5	17	00	12	16	39	17	07	12	30	30	18	01	11	45	3	18	07	11	11	19
16	11	00	3	3	17	00	13	13	8	17	07	13	34	34	18	01	12	10	2	18	07	12	11	19
16	11	01	48	48	17	00	14	10	1	17	07	14	54	78	18	01	13	24	24	18	07	13	11	19
16	11	02	144	14	17	00	15	11	1	17	07	15	87	14	18	01	14	13	2	18	07	14	11	19
16	11	03	3	11	17	00	16	13	8	17	07	16	45	44	18	01	15	2	2	18	07	15	11	19
16	11	04	164	11	17	00	17	11	1	17	07	17	38	38	18	01	16	9	9	18	07	16	11	19
16	11	05	3	14	17	00	18	13	11	17	07	18	54	44	18	01	17	2	2	18	07	17	11	19
16	11	06	3	11	17	00	19	14	8	17	07	19	35	35	18	01	18	5	5	18	07	18	11	19
16	11	07	3	11	17	00	20	14	8	17	07	20	33	33	18	01	19	3	3	18	07	19	11	19
16	11	08	3	11	17	00	21	11	1	17	07	21	33	33	18	01	20	7	7	18	07	20	11	19
16	11	09	3	11	17	00	22	11	1	17	07	22	36	36	18	01	21	3	3	18	07	21	11	19
16	11	10	3	11	17	00	23	11	1	17	07	23	36	36	18	01	22	3	3	18	07	22	11	19
16	11	11	3	11	17	00	24	11	1	17	07	24	36	36	18	01	23	3	3	18	07	23	11	19
16	11	12	3	11	17	00	25	11	1	17	07	25	36	36	18	01	24	3	3	18	07	24	11	19
16	11	13	3	11	17	00	26	11	1	17	07	26	36	36	18	01	25	3	3	18	07	25	11	19
16	11	14	3	11	17	00	27	11	1	17	07	27	36	36	18	01	26	3	3	18	07	26	11	19
16	11	15	3	11	17	00	28	11	1	17	07	28	36	36	18	01	27	3	3	18	07	27	11	19
16	11	16	3	11	17	00	29	11	1	17	07	29	36	36	18	01	28	3	3	18	07	28	11	19
16	11	17	3	11	17	00	30	11	1	17	07	30	36	36	18	01	29	3	3	18	07	29	11	19
16	11	18	3	11	17	00	31	11	1	17	07	31	36	36	18	01	30	3	3	18	07	30	11	19
16	11	19	3	11	17	00	32	11	1	17	07	32	36	36	18	01	31	3	3	18	07	31	11	19
16	11	20	3	11	17	00	33	11	1	17	07	33	36	36	18	01	32	3	3	18	07	32	11	19
16	11	21	3	11	17	00	34	11	1	17	07	34	36	36	18	01	33	3	3	18	07	33	11	19
16	11	22	3	11	17	00	35	11	1	17	07	35	36	36	18	01	34	3	3	18	07	34	11	19
16	11	23	3	11	17	00	36	11	1	17	07	36	36	36	18	01	35	3	3	18	07	35	11	19
16	11	24	3	11	17	00	37	11	1	17	07	37	36	36	18	01	36	3	3	18	07	36	11	19
16	11	25	3	11	17	00	38	11	1	17	07	38	36	36	18	01	37	3	3	18	07	37	11	19
16	11	26	3	11	17	00	39	11	1	17	07	39	36	36	18	01	38	3	3	18	07	38	11	19
16	11	27	3	11	17	00	40	11	1	17	07	40	36	36	18	01	39	3	3	18	07	39	11	19
16	11	28	3	11	17	00	41	11	1	17	07	41	36	36	18	01	40	3	3	18	07	40	11	19
16	11	29	3	11	17	00	42	11	1	17	07	42	36	36	18	01	41	3	3	18	07	41	11	19
16	11	30	3	11	17	00	43	11	1	17	07	43	36	36	18	01	42	3	3	18	07	42	11	19
16	11	31	3	11	17	00	44	11	1	17	07	44	36	36	18	01	43	3	3	18	07	43	11	19
16	11	32	3	11	17	00	45	11	1	17	07	45	36	36	18	01	44	3	3	18	07	44	11	19
16	11	33	3	11	17	00	46	11	1	17	07	46	36	36	18	01	45	3	3	18	07	45	11	19
16	11	34	3	11	17	00	47	11	1	17	07	47	36	36	18	01	46	3	3	18	07	46	11	19
16	11	35	3	11	17	00	48	11	1	17	07	48	36	36	18	01	47	3	3	18	07	47	11	19
16	11	36	3	11	17	00	49	11	1	17	07	49	36	36	18	01	48	3	3	18	07	48	11	19
16	11	37	3	11	17	00	50	11	1	17	07	50	36	36	18	01	49	3	3	18	07	49	11	19
16	11	38	3	11	17	00	51	11	1	17	07	51	36	36	18	01	50	3	3	18	07	50	11	19
16	11	39	3	11	17	00	52	11	1	17	07	52	36	36	18	01	51	3	3	18	07	51	11	19
16	11	40	3	11	17	00	53	11	1	17	07	53	36	36	18	01	52	3	3	18	07	52	11	19
16	11	41	3	11	17	00	54	11	1	17	07	54	36	36	18	01	53	3	3	18	07	53	11	19
16	11	42	3	11	17	00	55	11	1	17	07	55	36	36	18	01	54	3	3	18	07	54	11	19
16	11	43	3	11	17	00	56	11	1	17	07	56	36	36	18	01	55	3	3	18	07	55	11	19
16	11	44	3	11	17	00	57	11	1	17	07	57	36	36	18	01	56	3	3	18	07	56	11	19
16	11	45	3	11	17	00	58	11	1	17	07	58	36	36	18	01	57	3	3	18	07	57	11	19
16	11	46	3	11	17	00	59	11	1	17	07	59	36	36	18	01	58	3	3	18	07	58	11	19
16	11	47	3	11	17	00	60	11	1	17	07	60	36	36	18	01	59	3	3	18	07	59	11	19
16	11	48	3	11	17	00	61	11	1	17	07	61	36	36	18	01	60	3	3	18	07	60	11	19
16	11	49	3	11	17	00	62	11	1	17	07	62	36	36	18	01	61	3	3	18	07	61	11	19
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16	11	51	3	11	17	00	64	11	1	17	07	64	36	36	18	01	63	3	3	18	07	63	11	19
16	11	52	3	11	17	00	65	11	1	17	07	65	36	36	18	01	64	3	3	18	07	64	11	19
16	11	53	3	11	17	00	66	11	1	17	07	66	36	36	18	01	65	3	3	18	07	65	11	19
16	11	54	3	11	17	00	67	11	1	17	07	67	36	36	18	01	66	3	3	18	07	66	11	19
16	11	55	3	11	17	00	68	11	1	17	07	68	36	36										

factory agreement was soon found. All the  $F(hk1)$  were then calculated and their signs were used with the observed amplitudes in the calculation of a generalized projection (Cochran & Dyer, 1952) of the electron density on (001). This comprised calculation of  $C_1(XY)$ ,  $S_1(XY)$  and  $\sigma_1(XY)$  where

$$\sigma_1(XY) = c \int_0^1 \rho(XYZ) \cdot \exp 2\pi iZ \cdot dZ = C_1(XY) + iS_1(XY).$$

These projections showed the atoms clearly resolved and a single stage of refinement of the  $z$  parameters (from the peak heights in the various syntheses) gave the maps shown in Fig. 2. The  $z$ -coordinates were further refined by calculation of the  $hk2$  structure factors and generalized projections.

### Refinement of the structure

The structure has been refined in three cycles using the full three-dimensional data. The atomic coordinates were refined by the differential synthesis method and the best isotropic temperature factors were estimated by comparison of the 'observed' and 'calculated' peak heights and curvatures. The necessary computations were carried out on FERUT, the electronic digital computer at the University of Toronto, now known as TRANSFER at the Structures Laboratory of the National Research Council, Ottawa, using programmes which are readily applicable to all space groups and which take advantage of the symmetry elements of the space group under consideration (Ahmed & Barnes, 1958).

The first set of three-dimensional structure factors was calculated treating all the atoms, excluding hydrogen, as carbon and using the McWeeny (1951) carbon scattering-factor curve with  $B=3.93 \text{ \AA}$  for them all. The residual error index,  $R$ , at this stage was 0.25. As a result of the first differential synthesis using the observed structure amplitudes the atomic coordinates changed by an average of 0.033  $\text{\AA}$  and a maximum of 0.109  $\text{\AA}$ . The observed peak heights of the atoms at positions 5 and 19 were 7.76 and 7.45  $\text{e.\AA}^{-3}$  respectively while the heights at positions 10 and 24 were 9.70 and 9.33  $\text{e.\AA}^{-3}$ . From this stage, therefore, the latter two were taken to be the nitrogen positions and the McWeeny nitrogen  $f$ -curve was used for them.

The second set of structure factors was calculated with the atomic coordinates derived from the first observed differential synthesis and with five different isotropic temperature factors which varied from 3.60 to 4.40  $\text{\AA}$  depending on the distances of the atoms from their molecular centres. The two molecules were assumed to have the same thermal motion. The  $R$ -index was reduced by these calculations to 0.175 and the signs of 120 structure factors were changed. A second observed differential synthesis then gave an average shift of 0.005  $\text{\AA}$  (maximum 0.021  $\text{\AA}$ ) in the

atomic positions. A differential synthesis calculated from the calculated structure factors at this stage showed that the average error in an atomic coordinate due to termination of the series errors was 0.005  $\text{\AA}$  with a maximum 0.015  $\text{\AA}$ . Refinement of the isotropic temperature factors, from comparison of peak heights and curvatures in the 'observed' and 'calculated' syntheses, confirmed that the atoms further from the centre of each molecule have the higher thermal motion and indicated also that molecule  $B$  (atoms 15-28) as a whole is subject to more intense apparent thermal motion than is molecule  $A$ .

The atomic coordinates were corrected for finite summation errors by Booth's (1946) back-shift method and a third set of structure factors was calculated using different isotropic temperature factors for the individual atoms as shown in Table 2. The resultant  $R$ -index was 0.16.

The contributions to the structure factors of the 18 hydrogen atoms in the asymmetric unit, assumed to be situated 1.08  $\text{\AA}$  diagonally from the carbon positions with  $B=4.52 \text{ \AA}^2$  were next calculated, and added to the third set of calculated structure factors. Consequently  $R$  was reduced to 0.146, 14 reflexions changed sign and the signs of 52 reflexions became doubtful. Corrections were then made to the second observed differential synthesis which took account of the changed contributions of the 14 reflexions which had changed sign and eliminated the contributions of

Table 2. *Final atomic co-ordinates and isotropic temperature factors*

Molecule $A$	$x/a$	$y/b$	$z/c$	$B (\text{\AA}^2)$
$C_1$	-0.0262	0.1895	-0.2235	4.23
$C_4$	0.0422	0.2083	0.2126	
$C_6$	0.1743	-0.0279	0.1610	
$C_9$	0.0993	-0.0369	-0.2748	
$C_2$	-0.0419	0.2488	-0.1107	4.83
$C_3$	-0.0069	0.2588	0.1113	
$C_7$	0.1864	-0.0863	0.0316	
$C_8$	0.1480	-0.0903	-0.1866	
$C_{11}$	0.0257	0.1347	-0.1244	3.52
$C_{12}$	0.0597	0.1443	0.0990	
$C_{13}$	0.1239	0.0295	0.0727	
$C_{14}$	0.0867	0.0248	-0.1488	
$C_5$	0.1091	0.0896	0.1961	3.66
$N_{10}$	0.0386	0.0769	-0.2447	3.70
Molecule $B$				
$C_{15}$	0.2948	0.2139	0.3765	4.52
$C_{18}$	0.2421	0.1594	0.7718	
$C_{20}$	0.3745	-0.0686	0.5982	
$C_{23}$	0.4217	-0.0049	0.2120	
$C_{16}$	0.2490	0.2542	0.5071	4.99
$C_{17}$	0.2232	0.2265	0.7082	
$C_{21}$	0.4156	-0.1065	0.4548	
$C_{22}$	0.4404	-0.0742	0.2569	
$C_{25}$	0.3155	0.1421	0.4360	3.76
$C_{26}$	0.2891	0.1152	0.6393	
$C_{27}$	0.3528	0.0044	0.5561	
$C_{28}$	0.3767	0.0362	0.3571	
$C_{19}$	0.3082	0.0454	0.6956	3.97
$N_{24}$	0.3589	0.1038	0.3004	3.97



Table 3. *Final observed and calculated peak heights and curvatures*

Molecule <i>A</i>	Heights (e. Å <sup>-3</sup> )		Curvatures (e. Å <sup>-5</sup> )					
	<i>Q</i> <sub>o</sub>	<i>Q</i> <sub>c</sub>	$-A_{hh}^o$	$-A_{hh}^c$	$-A_{kk}^o$	$-A_{kk}^c$	$-A_{ll}^o$	$-A_{ll}^c$
C <sub>1</sub>	7.21	7.38	61.2	60.3	54.0	55.9	53.7	58.6
C <sub>4</sub>	6.96	7.17	58.6	57.2	48.9	54.1	52.7	56.2
C <sub>6</sub>	7.15	7.31	63.8	59.9	50.6	55.3	54.6	58.3
C <sub>9</sub>	7.54	7.51	65.8	62.3	59.3	59.6	58.5	60.6
C <sub>2</sub>	6.49	6.58	52.1	48.3	50.7	49.7	40.9	47.1
C <sub>3</sub>	6.26	6.63	48.9	50.0	46.9	49.7	40.5	48.2
C <sub>7</sub>	6.40	6.64	54.2	52.3	48.5	50.7	36.6	45.1
C <sub>8</sub>	6.70	6.84	56.6	54.6	52.8	51.4	44.2	51.9
C <sub>11</sub>	8.46	8.32	77.4	72.6	67.5	67.7	68.9	69.6
C <sub>12</sub>	8.20	8.21	75.9	71.5	63.5	68.1	65.2	64.8
C <sub>13</sub>	8.41	8.37	80.4	75.0	64.7	67.9	67.0	69.2
C <sub>14</sub>	8.42	8.38	79.7	74.3	67.0	68.7	68.1	69.9
C <sub>5</sub>	7.90	7.93	71.0	67.2	53.5	59.6	68.8	68.0
N <sub>10</sub>	9.92	9.89	84.7	81.4	75.9	78.3	83.6	81.2
Molecule <i>B</i>								
C <sub>15</sub>	7.10	7.23	57.6	56.5	59.0	57.5	54.9	58.4
C <sub>18</sub>	6.91	6.95	56.4	54.0	48.3	50.2	55.5	56.2
C <sub>20</sub>	6.61	7.06	46.8	52.5	54.3	54.8	49.0	56.1
C <sub>23</sub>	6.95	7.01	58.1	56.2	53.8	52.9	51.0	53.5
C <sub>16</sub>	6.39	6.48	50.9	48.8	49.2	46.6	42.0	46.1
C <sub>17</sub>	6.32	6.57	50.6	49.6	45.7	49.2	43.1	48.4
C <sub>21</sub>	6.07	6.42	45.5	48.6	49.0	48.5	35.0	43.3
C <sub>22</sub>	6.23	6.46	48.1	49.1	48.2	47.4	39.4	45.7
C <sub>25</sub>	7.99	8.03	72.7	69.2	62.9	62.3	65.7	67.3
C <sub>26</sub>	7.99	7.99	69.9	67.2	62.8	63.5	64.1	66.3
C <sub>27</sub>	7.74	7.89	66.5	65.9	60.0	61.7	60.9	63.7
C <sub>28</sub>	8.01	7.99	71.5	68.1	68.1	66.2	60.5	61.7
C <sub>19</sub>	7.65	7.78	63.4	63.9	61.5	63.3	63.8	64.3
N <sub>24</sub>	9.47	9.45	78.7	76.5	75.0	75.0	74.1	74.1

the 52 reflexions whose signs were now doubtfully determined. The mean change in the observed coordinates which resulted was 0.001 Å with maximum 0.004 Å. A second 'calculated' differential synthesis (including the *H*-contributions and eliminating the 52 reflexions of doubtful sign) was then computed to give final corrections for finite summation errors. These differed by an average of 0.001<sub>6</sub> Å from those indicated by the first 'calculated' differential synthesis. No attempt has been made to locate the hydrogen atoms in the structure more accurately.

The final atomic coordinates for the carbon and nitrogen atoms, after correction for finite summation errors, are given in Table 2 together with the individual temperature factors. The final observed and calculated peak heights and curvatures are given in Table 3. In these two tables those atoms which have been assigned the same isotropic temperature factors have been grouped together.

#### Molecular thermal motion

Comparison of the observed and calculated peak heights and curvatures shown in Table 3 suggests that the chosen isotropic temperature factors are not far from their best values and further that the chosen absolute scale of the measurements is very nearly correct. Detailed consideration of the ratios of the principal curvatures along *x*, *y* and *z* separately shows,

however, that more serious corrections to the temperature factors are required in order to allow for the anisotropic thermal motion of the individual atoms. These corrections have not been made in the present refinement but the nature of the thermal motion is nevertheless apparent in the results which have been obtained.

Higgs (1955) has shown that the mean-square amplitude of vibration in a molecular crystal is very largely determined by the rigid-body vibrations of the molecule as a whole and that the effects of these rigid-body vibrations can be resolved into components which arise from translation and restricted rotation (libration) of the molecule. Thus it appears that

$$\overline{u_i^2} \simeq P + QR_i^2,$$

where  $R_i$  is the distance of the atom from the molecular centre. Values of  $\overline{u_i^2} = 3B_i/8\pi^2$  for the different atomic groups in the two molecules of acridine II are plotted in Fig. 3 against  $R_i^2/L^2$  on the assumption that the molecules are formed from regular hexagons of side  $L$ . Corresponding results for acridine III and anthracene (cf. Part II) are plotted for comparison. It is clear that the apparent molecular thermal motions in the three structures are of the same type but that the translational components of the mean-square amplitudes of vibration appear to be higher in the acridines than in anthracene. This result is consistent

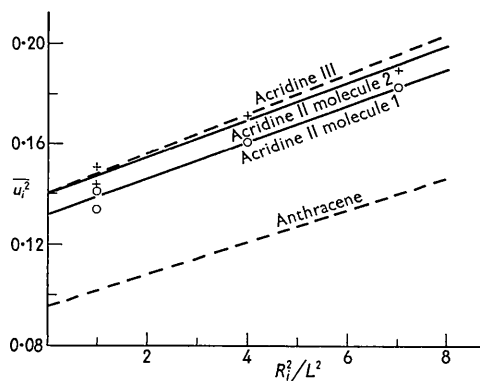


Fig. 3. Graphs of  $\overline{u_i^2}$  against  $R_i^2/L$  for atomic groups in acridine II, acridine III and anthracene.

with the lower melting-points of the acridines but it may be due in part at least to different systematic errors in the intensity data from which the structures have been determined and to a greater degree of static disorder in the acridines. In acridine II it appears that molecule *B* is subject to greater translational thermal motion (or disorder) than molecule *A* but that the librational motions of the two are very closely the same.

### Description of the structure

#### Arrangement of molecules

Fig. 4(a) and (b) show the arrangement of the molecules in the unit cell. As in acridine III the molecules

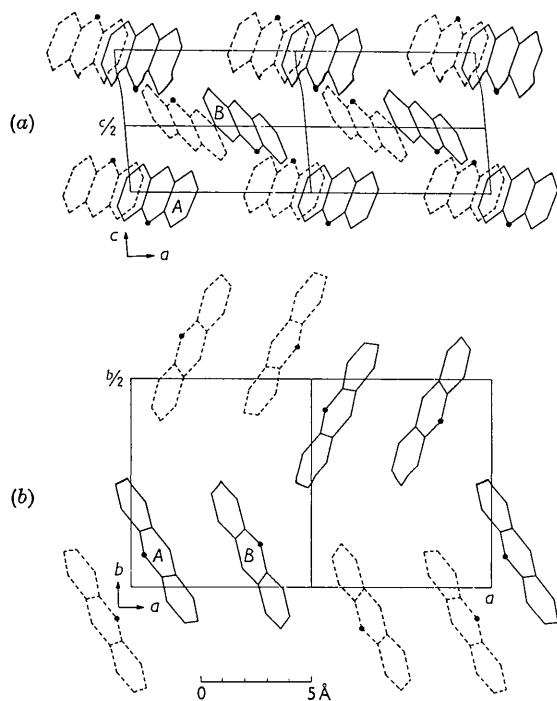


Fig. 4. Projections of the structure on (010) and (001). The two independent molecules in the asymmetric unit are labelled *A* and *B*. Nitrogen atoms are shown as full circles.

are arranged in anti-parallel pairs about the centres of symmetry but in this structure with its two molecules per asymmetric unit the arrangement is more complex. Molecule *A* and its symmetry-related fellows are arranged in pairs about the centres of symmetry at  $0, 0, 0$  and  $\frac{1}{2}, \frac{1}{2}, 0$  while the molecules *B* are arranged about the centres at  $\frac{1}{2}, 0, \frac{1}{2}$  and  $0, \frac{1}{2}, \frac{1}{2}$ . The structure thus may be described as being built up on (001) from layers of molecules in alternately the *A* and *B* orientations with the molecules in each layer fitting into the spaces in the layer below in a form of close-packing.

The packing units, each composed of two molecules, are not all equivalent. The anti-parallel molecules in orientation *A* are  $3.49 \text{ \AA}$  apart with minimum distance between atoms in adjacent molecules,  $C_1-C'_7$  and  $C_{12}-C'_9$ , equal to  $3.53 \text{ \AA}$ . (The hydrogen atoms have not been considered.) In normal projection one on the other the molecules in these pairs do not overlap exactly but are staggered in a very similar way to that found in acridine III and shown in Fig. 6(a) of Part II. The molecules in orientation *B*, on the other hand, are further apart than those in orientation *A*, interplanar spacing  $3.61 \text{ \AA}$ , and in normal projection one on the other only the terminal rings overlap. The shortest distance between atoms in adjacent molecules in these pairs is  $3.69 \text{ \AA}$  between  $C_{20}$  and  $C'_{25}$ .

These differences in the packing of the two molecules in the asymmetric unit are consistent with the observed differences in their apparent thermal motions. The fact that molecules *A* which are the more intimately related in pairs appear to be subject to less intense thermal motion than molecules *B* may indicate that molecules *A* have the greater tendency to vibrate as two-fold molecular units. It must be remembered, however, that the observations could be explained also by a greater degree of disorder in the packing of molecules *B*.

The close packing of the two-fold molecular groups depends principally on intermolecular contacts of the kind usually found between benzene rings: all the intermolecular C-C distances concerned are larger than  $3.6 \text{ \AA}$ . The N-C<sub>5</sub> distance between molecules related by the *c*-axis unit-cell translation is  $3.69 \text{ \AA}$  for molecules *A* and  $3.85 \text{ \AA}$  for molecules *B*.

A preliminary analysis of this packing in relation to that found in anthracene has been given by Miasnikova & Kitaigorodskii (1958) in discussion of their investigation of mixed crystals of the acridine-anthracene system. They have shown that the addition of only 7% anthracene prevents the growth of mixed crystals of the acridine II type and, making use of the results presented here which were communicated to them privately, they have shown that this effect arises from the fact that the acridine molecules in this structure pack together more closely than anthracene molecules are able to do. A more detailed comparison of the structures of acridine II, acridine III and anthracene, in relation to the polymorphism of acridine, will be given elsewhere.

*Molecular dimensions and accuracy*

Bond lengths and angles within the individual molecules, calculated from the final atomic coordinates, are shown in Fig. 5(a), (b), (c) and (d).

The factors affecting the accuracy of structure determinations such as this in which the refinement is incomplete were discussed in detail in Part II. They must be taken into account in assessing the accuracy of the present results. Here again the standard deviation of the atomic coordinates have been calculated using the residuals  $\Delta F = F_o - F_c$  as estimates of  $\sigma(F)$ . Their root-mean-square values vary between 0.0037 Å for N<sub>10</sub> and 0.0072 Å for C<sub>21</sub>. The standard deviations of the bond lengths and angles, estimated by the use of Ahmed & Cruickshank's (1952) formulae, vary between 0.006 and 0.011 Å for lengths and 20' and 32' for angles.

The differences  $\Delta(l)$  and  $\Delta(\theta)$  between chemically equivalent bonds and angles within the two molecules are shown in Fig. 5(e) and (f). Estimation of the Student's distribution,  $t_0 = (l_1 - l_2)/(S_1^2 - S_2^2)^{1/2}$  (Cruickshank & Robertson, 1953) for these differences indicates that:

(i) for both molecules the differences in the bond lengths are not significant, even for the largest dif-

ference 0.021 Å between the bonds C<sub>3</sub>-C<sub>4</sub> and C<sub>6</sub>-C<sub>7</sub> for which  $t_0 = 1.81$  and  $P > 0.05$ ;

(ii) for molecule *A* the differences in the angles are not significant except for the pair of angles C<sub>1</sub>-C<sub>11</sub>-C<sub>12</sub> and C<sub>9</sub>-C<sub>14</sub>-C<sub>13</sub> for which the difference of 62' is possibly significant ( $t_0 = 2.09$ ;  $0.05 > P > 0.01$ );

(iii) for molecule *B* the differences in the angles are not significant for 6 pairs, possibly significant for two pairs with  $\Delta(\theta) = 75'$  and  $78'$ , and significant for the pair C<sub>18</sub>-C<sub>26</sub>-C<sub>25</sub> and C<sub>20</sub>-C<sub>27</sub>-C<sub>28</sub> for which  $\Delta(\theta) = 101'$ ,  $t_0 = 2.97$ ,  $0.01 > P > 0.001$ .

Thus there are indications that the chemically equivalent halves of the two molecules may be slightly different in the crystal structure and it is worth remarking that the most significant differences are between bond-angles which may be more readily affected by molecular interactions than bond lengths. The differences are not well established, however, and in view of the uncertainty with which the standard deviations are estimated and the presence of systematic effects not taken into account in the  $F_c$  they may be a better indication of the errors remaining in the results than of real differences within the molecules. The weighted means of chemically equivalent bond lengths

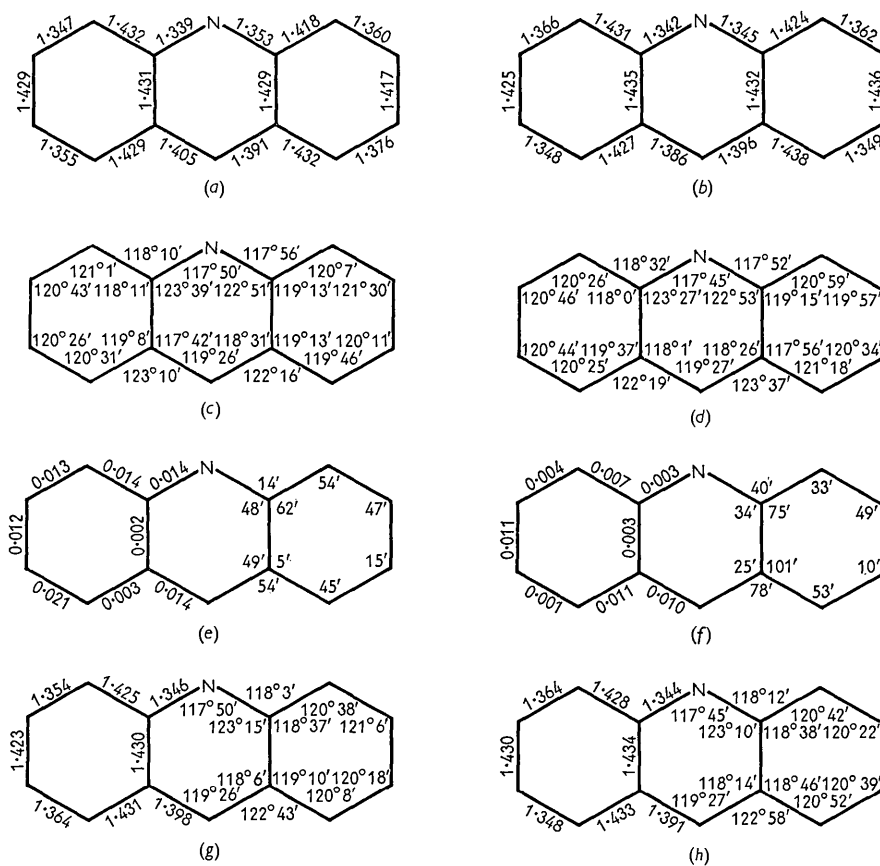


Fig. 5. Bond lengths (Å) and angles in the two molecules *A* and *B*. (a, b) observed bond lengths; (c, d) observed bond angles; (e, f) differences between chemically equivalent bond lengths and angles; (g, h) weighted mean bond lengths and angles.

Table 4. Mean bond lengths in acridine II and III, anthracene and phenazine

Bond	Acridine II	$\sigma(l)$	Acridine III	Anthracene	Phenazine	With correction for librations	
						Acridine II	Anthracene
C <sub>1</sub> -C <sub>2</sub>	1.358	0.004	1.372	1.364	1.375	1.360	1.366
C <sub>3</sub> -C <sub>4</sub>	1.357	0.004	1.370	1.364	1.375	1.359	1.366
C <sub>1</sub> -C <sub>11</sub>	1.427	0.003	1.431	1.431	1.408	1.429	1.433
C <sub>4</sub> -C <sub>12</sub>	1.432	0.004	1.425	1.431	1.408	1.434	1.433
C <sub>2</sub> -C <sub>3</sub>	1.426	0.005	1.417	1.417	1.412	1.428	1.419
C <sub>11</sub> -C <sub>12</sub>	1.432	0.003	1.432	1.434	1.433	1.434	1.436
C <sub>5</sub> -C <sub>12</sub>	1.395	0.004	1.389	1.396	—	1.398	1.399
C <sub>11</sub> -N <sub>10</sub>	1.345	0.003	1.344	—	1.345	1.348	—

and angles within the two molecules separately have therefore been calculated for comparison of the two molecules on the assumption that the observed differences are not real. These bond lengths and angles are shown in Fig. 5(g) and (h).

The differences in bond lengths between the two molecules are all not significant, including the difference of 0.016 Å between C<sub>3</sub>-C<sub>4</sub> and C<sub>17</sub>-C<sub>18</sub> for which  $t_0=1.41$ ,  $P>0.05$ . It seems reasonable to assume, therefore, at least for the bond lengths, that there are no real differences between the two molecules and that the four observed values for each bond can be averaged for comparison with similar measurements in acridine III and anthracene. These values are shown in Table 4 together with their standard deviations estimated on the assumption that each mean value is derived from four independent measurements of the same quantity. Corresponding results for acridine III, anthracene and phenazine (Herbstein & Schmidt, 1955) are also shown.

These bond lengths are subject to systematic errors due to the librations of the molecules (Cox, Cruickshank & Smith, 1955). Corrections for these errors have been applied in the latest refinement of the structure of anthracene (Cruickshank, 1956) and in view of the fact that the librational motion of the molecules in the acridines has been shown to be closely similar to that in anthracene it is reasonable to suppose that the corrections required here are the same. The bond angles are not much affected by these errors but intramolecular bond lengths are too short by 0.002 or 0.003 Å. The corrected bond lengths are shown in the last columns of Table 4.

### Molecular planarity

The molecules are not strictly planar. The best planes through the two molecules *A* and *B* are given by

$$0.8359x' + 0.4340y - 0.3359z' = 1.7465 \quad (1)$$

and

$$0.8144x' + 0.3056y + 0.4933z' = 6.0994 \quad (2)$$

respectively, where  $x'=x+z \cos \beta$ ,  $z'=z \sin \beta$  are the atomic coordinates measured in Ångström units in the orthogonal coordinate system  $a, b, c^*$ . The perpendicular distances of the atoms of each molecule from the

mean plane of that molecule are shown in Fig. 6(a) and (b). Application of the  $\chi^2$  test shows that for molecule *A* as a whole  $\chi^2=53$ ,  $\nu=11$ ,  $P<0.0001$  and for molecule *B*,  $\chi^2=272$ ,  $\nu=11$ ,  $P<0.0001$ . The deviations of both molecules from planarity thus are 'highly significant'.

In the search for systematic features in these deviations from planarity it is illuminating to consider the deviations of the atoms from the best planes through the central pyridine rings alone. These planes for the molecules *A* and *B* are

$$0.8392x' + 0.4338y - 0.3280z' = 1.7526 \quad (3)$$

and

$$0.8153x' + 0.3057y + 0.4916z' = 6.1211 \quad (4)$$

respectively, and the perpendicular distances of the atoms from them are shown in Figs. 6(c) and (d). For the central ring of molecule *A*,  $\chi^2=4.6$ ,  $\nu=3$ ,  $P>0.05$  and for the central ring of molecule *B*,  $\chi^2=6.7$ ,  $\nu=3$ ,  $P>0.05$ . The central rings of the two molecules thus do not deviate significantly from planarity. Inspection of the deviations shown in Figs. 6(c) and (d) suggests, however, that the distortion of molecule *A* may be due to a tendency of the central ring to take up a 'chair' configuration to allow for inclusion of the nitrogen atom without upsetting the planar and symmetrical configuration of bonds at the carbon atoms. This is suggested by the fact that the bonds C<sub>11</sub>-C<sub>1</sub> and C<sub>14</sub>-C<sub>9</sub> appear to be directed down below the plane of the diagram (and towards the centre of symmetry at 0, 0, 0 and the symmetry related molecule) while the bonds C<sub>12</sub>-C<sub>4</sub> and C<sub>13</sub>-C<sub>6</sub> appear to be directed above the plane of the diagram. The best planes of the two outer rings are roughly parallel to one another but are rotated about the length of the molecule with respect to the best plane of the central ring. In molecule *B* on the other hand the central ring appears to have taken up a 'boat' configuration so that the bonds C<sub>25</sub>-C<sub>15</sub>, C<sub>26</sub>-C<sub>18</sub>, C<sub>27</sub>-C<sub>20</sub> and C<sub>28</sub>-C<sub>23</sub> are all directed down below the plane of the diagram and away from the centre of symmetry at  $\frac{1}{2}, 0, \frac{1}{2}$ .

Even when uncertainties in estimation of the standard deviations are taken into account it seems certain that molecule *B* is not planar. The distortion is described fairly well by bending across the line N-C<sub>19</sub> through about 2° and it is particularly interest-

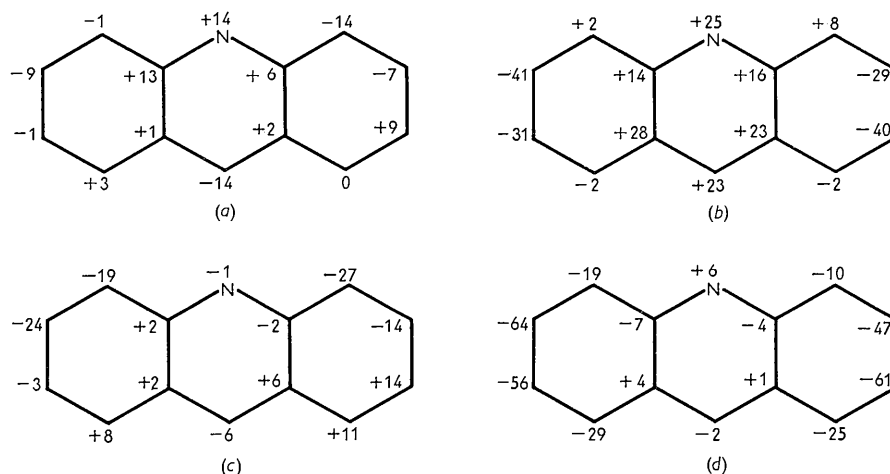


Fig. 6. Perpendicular distances of atoms from possible molecular planes: (a) molecule *A* and plane 1; (b) molecule *B* and plane 2; (c) molecule *A* and plane 3; (d) molecule *B* and plane 4. All distances are in  $\text{\AA} \times 10^3$  and positive signs indicate increased distance from the origin.

ing that this result should be so similar to that obtained in acridine III. The deviations from planarity of molecule *A* are not so certainly significant, but interpreted in the way outlined above they too appear to be systematic and meaningful.

### Discussion

The results of this investigation are in good agreement with the results obtained earlier by a study of acridine III. In particular the two independent structure determinations have given consistent estimations of the molecular dimensions and also consistent indications that molecules of acridine in the crystalline state are distorted from the ideally planar form. The implications of these results have been discussed already in Part II and are not considered further here.

It is evident that closer investigation of the effects of molecular interactions in acridine II depends first on the collection of more accurate intensity data and then on complete refinement of the structure to fit those data. The results presented here show that such an investigation may give interesting results.

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